

Abstract Algebra

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1. GROUPS

DEFINITION. A group is a set G with a law of composition that has the following properties:

- Closure (Totality): $\forall a, b \in G, ab \in G$

- *Associativity:* $\forall a, b, c \in G, (ab)c = a(bc)$
- *Identity:* $\exists 1 \in G, 1a = a1 = a$
- *Invertibility:* $\exists b \in G, ba = ab = 1$

A group is abelian or commutative if it satisfies $\forall a, b \in G, ab = ba$.

DEFINITION. A subset H of a group G is called a subgroup if it has the following properties:

- *Closure (Totality):* $\forall a, b \in H, ab \in H$
- *Identity:* $1 \in H$
- *Inverses:* If $a \in H, a^{-1} \in H$

1.1. Homomorphisms

Let G, G' be groups. A homomorphism $\varphi: G \rightarrow G'$ is any map satisfying the rule

$$\varphi(ab) = \varphi(a)\varphi(b) \quad (1)$$

for all $a, b \in G$. When the map φ is bijective, it's called an isomorphism.

1.2. Conjugation

The conjugation is a map $G \times G \rightarrow G$, defined by $\{gxg^{-1}, g, x \in G\}$.

DEFINITION. The stabilizer of an element $x \in G$ for the operation of conjugation is called the centralizer of x and is denoted by $Z(x)$:

$$Z(x) = \{g \in G | gxg^{-1} = x\} = \{g \in G | gx = xg\} \quad (2)$$

Note $x \in Z(x)$, because x commutes with itself.

DEFINITION. The orbit of x for the operation of conjugation is called the conjugacy class of x .

$$C_x = \{x' \in G | x' = gxg^{-1} \text{ for some } g \in G\} \quad (3)$$

By definition, $|G| = |C_x||Z(x)|$.

DEFINITION. If H is a subgroup of G , and satisfies $gHg^{-1} = H$, which means H is self-conjugate for all elements $g \in G$, then H is called an invariant subgroup or a normal subgroup of G .

1.3. Vector Space

DEFINITION. A real vector space is a set V together with two laws of composition:

- Addition:* $V + V \rightarrow V$, written $v, w \rightsquigarrow v + w$
- Scalar multiplication:* $\mathbb{R} \times V \rightarrow V$, written $c, v \rightsquigarrow cv$

The laws of composition satisfies the following axioms:

- Addition makes V into an abelian group V^+ ;*
- Scalar multiplication is associative: $(ab)v = a(bv)$;*
- Identity operation: $1v = v$;*
- Two distributive laws hold: $(a + b)v = av + bv$ $a(v + w) = av + aw$.*

DEFINITION. The dimension of a finite-dimensional vector space V is the number of vectors in a basis. The dimension is denoted by $\dim V$.

Given a basis (v_1, v_2, \dots, v_n) of a vector space V , for all $v \in V$, it can be expressed as

$$v = x_1v_1 + \dots + x_nv_n \quad (4)$$

where the scalar x_i are called the coordinates of v .

DEFINITION. Suppose V is a vector space, and S is a subspace of V , then the span of S is

$$\text{Span } S = \{v | v = c_1 v_1 + \dots + c_r v_r\} \quad \text{where } v_1, \dots, v_r \in S \quad (5)$$

If W_1, W_2, \dots, W_n are subspaces of a vector space V , the span of the subspaces is denoted by

$$W_1 + \dots + W_n = \{v \in V | v = w_1 + \dots + w_n, \text{ with } w_i \in W_i\} \quad (6)$$

The sum is the smallest subspace containing W_1, W_2, \dots, W_n . The subspaces are independent if

$$w_1 + \dots + w_n = 0, \text{ with } w_i \in W_i \text{ implies } w_i = 0 \text{ for all } i. \quad (7)$$

DEFINITION. The direct sum of W_1, \dots, W_n is

$$V = W_1 \oplus W_2 \oplus \dots \oplus W_n, \text{ if } V = W_1 + W_2 + \dots + W_n \text{ and if } W_1, \dots, W_n \text{ are independent} \quad (8)$$

Let W_1, W_2 be subspaces of a finite-dimensional vector space V , then

$$\dim(W_1) + \dim(W_2) = \dim(W_1 \cap W_2) + \dim(W_1 + W_2) \quad (9)$$

DEFINITION. The Kronecker product of an $m \times n$ matrix A and an $p \times q$ matrix B is defined as

$$A \otimes B = \begin{pmatrix} a_{11}B & \dots & a_{1n}B \\ \vdots & \ddots & \vdots \\ a_{m1}B & \dots & a_{mn}B \end{pmatrix} \quad (10)$$

which is also called a direct product or a tensor product.

DEFINITION. A metric function on a vector space V is a mapping of a pair of vectors into a number in the field F associated with the vector space.

$$(v_1, v_2) = f \quad v_1, v_2 \in V, f \in F \quad (11)$$

This mapping obeys

$$(v_1, \alpha v_2 + \beta v_3) = \alpha(v_1, v_2) + \beta(v_1, v_3) \quad (12)$$

$$(\alpha v_1 + \beta v_2, v_3) = (v_1, v_3)\alpha + (v_2, v_3)\beta \quad (13)$$

$$(\alpha v_1 + \beta v_2, v_3) = (v_1, v_3)\alpha^* + (v_2, v_3)\beta^* \quad (14)$$

Metrics obeying conditions (12) and (13) are called bilinear metrics; those obeying (12) and (14) are called sesquilinear.

DEFINITION. Groups preserving bilinear symmetric metrics are called orthogonal.

DEFINITION. Groups preserving bilinear antisymmetric metrics are called symplectic.

DEFINITION. Groups preserving sesquilinear symmetric metrics are called unitary.

1.4. Linear Groups

The real general linear group GL_n is [1]

$$GL_n = \{P \in GL_n(\mathbb{R}) | P^{-1} \in GL_n(\mathbb{R})\} \quad (15)$$

The real special linear group SL_n is

$$SL_n = \{P \in GL_n(\mathbb{R}) | \det P = 1\} \quad (16)$$

The orthogonal group O_n is

$$O_n = \{P \in GL_n(\mathbb{R}) | P^T P = I\} \quad (17)$$

The unitary group U_n is

$$U_n = \{P \in GL_n(\mathbb{C}) | P^*P = I\} \quad (18)$$

The symplectic group SP_{2n} is

$$S = \begin{pmatrix} 0 & I \\ -I & 0 \end{pmatrix}$$

$$SP_{2n} = \{P \in GL_{2n}(\mathbb{R}) | P^t S P = S\} \quad (19)$$

The orthogonal group for indefinite forms is

$$O_{3,1} = \{P \in GL_n(\mathbb{R}) | P^t I_{3,1} P = I_{3,1}\} \quad (20)$$

The linear operators represented by $O_{3,1}$ are called Lorentz transformations. The dimensionalities relations of classical groups is in Fig. 2.5 in [2].

AXIOM. If $A \in O_n$, then $|\lambda(A)| = 1$ and $\det(A) = \pm 1$.

Proof. Let λ be an eigenvalue of A and $Av = \lambda v$ with $v \neq 0$, then

$$\begin{aligned} (Av)^T(Av) &= (\lambda v)^T(\lambda v) \\ v^T(A^T A)v &= \lambda^2 v^T v \end{aligned}$$

Since $A \in O_n$ and $A^T A = I$,

$$\begin{aligned} v^T v &= \lambda^2 v^T v \\ \lambda^2 &= 1 \\ |\lambda| &= 1 \end{aligned} \quad (21)$$

It should be noted λ can be complex. As for the determinant, we have

$$\begin{aligned} \det(A^T A) &= \det(I) = 1 \\ (\det(A))^2 &= 1 \\ \det(A) &= \pm 1 \end{aligned}$$

When $\det(A) = 1$, it represents rotation; otherwise $\det(A) = -1$, it represents reflection. \square

THEOREM. (EULER'S THEOREM) If $A \in SO_3$, then at least one eigenvalue of A is 1.

Proof. Let λ be an eigenvalue of A and $Av = \lambda v$ with $v \neq 0$, then the characteristic equation is

$$f(\lambda) = \lambda^3 + a\lambda^2 + b\lambda + c = 0$$

Its solution satisfies

$$\begin{aligned} \lambda_1 \lambda_2 \lambda_3 &= \det(A) = 1 \\ \lambda_1 + \lambda_2 + \lambda_3 &= \text{trace}(A) \end{aligned}$$

Since the highest order of $f(\lambda)$ is an odd number, then it must have a real root λ_1 .

1. If three roots are all real, from (21), $\lambda = \pm 1$. Considering $\lambda_1 \lambda_2 \lambda_3 = 1$, three roots can't be all -1 . Therefore there is at least one root is 1.
2. If two roots are complex and one root λ_1 is real, then $\lambda_1 = \pm 1$ and

$$f(\lambda) = (\lambda - \lambda_1)(\lambda^2 + e\lambda + f) = 0$$

then $\lambda_{2,3} = \alpha \pm \beta i$ and $\lambda_2 \lambda_3 = \alpha^2 + \beta^2 > 0$. Since we have $\lambda_1 \lambda_2 \lambda_3 = 1 > 0$, it's true that $\lambda_1 > 0$. Moreover, $\lambda_1 = 1$.

Above all, at least one eigenvalue of A is 1. \square

1.5. Linear Transformations

Let $T: V \rightarrow W$ be any linear transformation. We introduce two subspaces

$$\begin{aligned} \text{kernel of } T: \quad \text{Ker } T &= \{v \in V | T(v) = 0\} \\ \text{image of } T: \quad \text{im } T &= \{w \in W | w = T(v) \forall v \in V\} \end{aligned}$$

Assuming V is finite-dimensional, then

$$\dim(V) = \dim(\text{ker } T) + \dim(\text{im } T) = \text{rank} + \text{nullity} \quad (22)$$

1.6. Special Unitary Group

The elements of SU_2 are complex 2×2 matrices of the form

$$P = \begin{pmatrix} a & b \\ -\bar{b} & \bar{a} \end{pmatrix}, \text{ with } \bar{a}a + \bar{b}b = 1 \quad (23)$$

Proof. Because $\det P = 1$, let

$$\begin{aligned} P &= \begin{pmatrix} a & b \\ u & v \end{pmatrix} \\ \begin{pmatrix} \bar{a} & \bar{u} \\ \bar{b} & \bar{v} \end{pmatrix} &= P^* = P^{-1} = \begin{pmatrix} v & -b \\ -u & a \end{pmatrix} \end{aligned}$$

Therefore

$$\begin{aligned} v &= \bar{a} \\ u &= -\bar{b} \\ \det P &= \bar{a}a + \bar{b}b = 1 \end{aligned}$$

□

Writing $a = x_0 + x_1i, b = x_2 + x_3i$, then $\det P = \bar{a}a + \bar{b}b = x_0^2 + x_1^2 + x_2^2 + x_3^2 = 1$, which defines a unit 3-sphere in \mathbb{R}^4 .

$$\begin{aligned} SU_2 &\longleftrightarrow \mathbb{S}^3 \\ P &= \begin{pmatrix} x_0 + x_1i & x_2 + x_3i \\ -x_2 + x_3i & x_0 - x_1i \end{pmatrix} \longleftrightarrow (x_0, x_1, x_2, x_3) \\ \mathbf{1} &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \longleftrightarrow (1, 0, 0, 0) \\ \mathbf{i} &= \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} \longleftrightarrow (0, 1, 0, 0) \\ \mathbf{j} &= \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \longleftrightarrow (0, 0, 1, 0) \\ \mathbf{k} &= \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \longleftrightarrow (0, 0, 0, 1) \end{aligned}$$

A famous theorem of topology asserts that the only spheres on which one can define continuous group laws are the 1-sphere and the 3-sphere.

1.7. Matrix Representation

An n -dimensional matrix representation of a group G is a homomorphism

$$\rho: G \rightarrow GL_n(V) \quad (24)$$

The dimension of the representation ρ is defined to be the dimension of the finite-dimensional vector space V .

If Γ is a mapping from A to B such that $\forall a \in A, \Gamma(a) \in B$, then Γ is said to be into, or injective. If $\forall b \in B, \exists a \in A, \Gamma(a) = b$, then Γ is said to be onto, or surjective.

1.8. Cosets

DEFINITION. The set denoted by aN is called a coset of N in G if:

$$aN = \{g \in G \mid g = an \text{ for some } n \in N\} \quad (25)$$

COROLLARY. A group homomorphism $\varphi: G \rightarrow G'$ is injective (an isomorphism) iff its kernel is the trivial subgroup.

2. FIELDS

DEFINITION. A field F is a set together with two laws of composition:

$$F + F \rightarrow F, F \times F \rightarrow F$$

called addition and multiplication, which satisfy these axioms:

- Addition: abelian group F^+ ; identity element 0
- Multiplication: commutative, abelian group F^\times ; identity element 1
- Distributive law: $\forall a, b, c \in F, a(b + c) = ab + ac$

DEFINITION. A vector space V over a field F is defined as in (1.3), with F replacing \mathbb{R} .

- a) Addition: $V + V \rightarrow V$, written $v, w \rightsquigarrow v + w$
- b) Scalar multiplication: $F \times V \rightarrow V$, written $c, v \rightsquigarrow cv$

The laws of composition satisfies the following axioms:

- i. Addition makes V into an abelian group V^+ ;
- ii. Scalar multiplication is associative: $(ab)v = a(bv), \quad \forall a, b \in F, \forall v \in V$;
- iii. Identity operation: $1v = v \quad \forall 1 \in F, \forall v \in V$;
- iv. Two distributive laws hold: $(a + b)v = av + bv \quad a(v + w) = av + aw \quad \forall a, b \in F, \forall v, w \in V$.

3. RINGS

DEFINITION. A ring R is a set together with two laws of composition $+$ and \times , called addition and multiplication, that satisfy these axioms:

- Addition: abelian group R^+ ; identity element 0
- Multiplication: commutative and associative, identity element 1
- Distributive law: $\forall a, b, c \in R, a(b + c) = ab + ac$

4. LIE GROUPS AND LIE ALGEBRAS

4.1. Lie Groups

DEFINITION. A topological space T is a set of points on which is placed a topology \mathcal{T} . The topology \mathcal{T} is a choice (set) of subsets S_1, S_2, \dots of T : $S_i \subset T, S_i \in \mathcal{T}$. The topology \mathcal{T} obeys

AXIOM 1. The empty set $\Phi \in \mathcal{T}$; the topological space $T \in \mathcal{T}$.

AXIOM 2. Finite intersections of elements in \mathcal{T} are elements in \mathcal{T} .

$$\bigcap_{i=1}^{\text{finite}} S_i \in \mathcal{T} \quad (26)$$

AXIOM 3. Arbitrary unions of elements in \mathcal{T} are elements in \mathcal{T} .

$$\bigcup_{i=1}^{\text{finite or infinite}} S_i \in \mathcal{T} \quad (27)$$

The elements S_i in the topology \mathcal{T} are called open sets.

A topological space obeying the additional axiom 4 is called a Hausdorff space.

AXIOM 4. If $p \in T, q \in T, p \neq q$, then there exist $S_p \in \mathcal{T}, S_q \in \mathcal{T}$ with the property $p \in S_p, q \in S_q, S_p \cap S_q = \Phi$.

DEFINITION. An open set S_p containing p is called a neighborhood of p , symbolically $p \in S_p \in \mathcal{T}$.

DEFINITION. A space T is compact if every infinite sequence of points $t_1, t_2, \dots, (t_i \in T)$ contains a subsequence of points that a) converges to a point and b) this point is in T . This point is called a limit point.

DEFINITION. A space T is closed if it contains all its limit points. The set T , together with all its limit points, is called the closure \bar{T} of T .

DEFINITION. $\phi: (T, \mathcal{T}) \rightarrow (U, \mathcal{U})$ is continuous if the inverse image of any open set in U is an open set in T .

DEFINITION. A differentiable manifold \mathcal{M} consists of

1. A Hausdorff space (T, \mathcal{T}) ;
2. A collection Φ of mappings $\phi_p \in \Phi: \phi_p: T \rightarrow R_\eta \quad p \in T$.

which obeys the following properties

- a) ϕ_p is a 1-1 mapping of an open set $T_p (p \in T_p)$ into an open set in R_η ;
- b) $\bigcup T_p = T$;
- c) If $T_p \cap T_q$ is not empty, then $\phi_p(T_p \cap T_q)$ is an open set in R_η , and $\phi_q(T_p \cap T_q)$ is an open set in R_η which is different from $\phi_p(T_p \cap T_q)$. The mapping $\phi_p \circ \phi_q^{-1}$ must be continuous and differentiable;
- d) (Maximality). The mappings $\phi_p \circ \phi_q^{-1}$ and $\phi_q \circ \phi_p^{-1}$, described in c), are mappings in Φ .

DEFINITION. A topological group or a continuous group consists of

1. An underlying η -dimensional manifold \mathcal{T} ;
2. An operation ϕ mapping each pair of points (β, α) in the manifold into another point γ in the manifold $\phi: \mathcal{T} \times \mathcal{T} \rightarrow \mathcal{T}$;
3. In terms of coordinate systems around the points γ, β, α , we write

$$\gamma^\mu = \phi^\mu(\beta^1, \dots, \beta^\eta; \alpha^1, \dots, \alpha^\eta); \quad \mu = 1, 2, \dots, \eta \quad (28)$$

The functions

$$\begin{aligned} \phi: \beta \times \alpha &\rightarrow \gamma = \beta\alpha \\ \psi: \alpha &\rightarrow \alpha^{-1} \end{aligned}$$

must be continuous and satisfy:

- a) Closure: $\gamma^\mu = \phi^\mu(\beta, \alpha) \quad \alpha, \beta, \gamma \in \mathcal{T}$;
- b) Associativity: $\phi^\mu(\gamma, \phi^\mu(\beta, \alpha)) = \phi^\mu(\phi^\mu(\gamma, \beta), \alpha)$;
- c) Identity: $\phi^\mu(e, \alpha) = \alpha^\mu = \phi^\mu(\alpha, e)$;
- d) Inverse: $\phi^\mu(\alpha^{-1}, \alpha) = e^\mu = \phi^\mu(\alpha, \alpha^{-1})$.

DEFINITION. A continuous group of transformations consists of

1. An underlying η -dimensional manifold \mathcal{T} ;
2. An operation ϕ mapping each pair of points (β, α) in the manifold into another point γ in the manifold $\phi: \mathcal{T} \times \mathcal{T} \rightarrow \mathcal{T}$, which obeys the postulates of a topological group;
3. A geometric space G , which is an N -dimensional manifold, and a mapping $f: \mathcal{T} \times G \rightarrow G$.

$$y^i = f^i(\alpha^1, \dots, \alpha^\eta; x^1, \dots, x^N); \quad i = 1, 2, \dots, N \quad (29)$$

The function must be continuous and obey

- a) Closure: $y^i = f^i(\alpha, x) \quad \alpha \in \mathcal{T}, x \in G, y \in G$;
- b) Associativity: $f^i(\beta, f^i(\alpha, x)) = f^i(\phi^\mu(\beta, \alpha), \alpha)$;
- c) Identity: $f^i(e, x) = x^i = f^i(x, e)$;
- d) Inverse: $f^i(\alpha^{-1}, f^i(\alpha, x)) = f^i(\alpha, f^i(\alpha^{-1}, x)) = f^i(\phi^\mu(\alpha, \alpha^{-1}), x) = x^i$.

Every continuous group may be considered as a continuous group of transformations if we allow it to act on itself: $G \equiv \mathcal{T} \quad f \equiv \phi$.

Example. Consider a coordinate transformation $f(\alpha^1, \alpha^2; x) = \alpha^1 x + \alpha^2$, where $\alpha^1 \neq 0$, f is defined on the geometric space $G = R_1$. The function ϕ is defined as

$$\begin{aligned} x'' &= f(\beta^1, \beta^2; x') \quad x' = f(\alpha^1, \alpha^2; x) \\ x'' &= \beta^1 x' + \beta^2 = \beta^1(\alpha^1 x + \alpha^2) + \beta^2 = \beta^1 \alpha^1 x + \beta^1 \alpha^2 + \beta^2 = f(\beta^1 \alpha^1, \beta^1 \alpha^2 + \beta^2; x) \\ \phi(\beta^1, \beta^2; \alpha^1, \alpha^2) &\equiv (\beta^1 \alpha^1, \beta^1 \alpha^2 + \beta^2) \end{aligned} \quad (30)$$

where ϕ acts on the topological space $\mathcal{T} = R_2$. The above definitions form a continuous group of transformation.

The nonsingular matrices given by

$$(\alpha^1, \alpha^2) \leftrightarrow \begin{pmatrix} \alpha^1 & \alpha^2 \\ 0 & 1 \end{pmatrix} \quad (31)$$

is a representation of this group in terms of 2×2 matrices, where

$$f: \begin{pmatrix} x' \\ 1 \end{pmatrix} = \begin{pmatrix} \alpha^1 & \alpha^2 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ 1 \end{pmatrix} = \begin{pmatrix} \alpha^1 x + \alpha^2 \\ 1 \end{pmatrix} \quad (32)$$

$$\phi: \begin{pmatrix} \gamma^1 & \gamma^2 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} \beta^1 & \beta^2 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \alpha^1 & \alpha^2 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} \beta^1 \alpha^1 & \beta^1 \alpha^2 + \beta^2 \\ 0 & 1 \end{pmatrix} \quad (33)$$

Additionally we have

$$(x')^N = (\alpha^1 x + \alpha^2)^N = \sum_{r=0}^N \binom{N}{r} (\alpha^1)^r (\alpha^2)^{N-r} x^r \quad (34)$$

The $N+1$ homogeneous polynomials $(x^N, x^{N-1}, \dots, x, 1)$ can be used as bases for an $(N+1) \times (N+1)$ matrix representation of the projective group. For $N=3$,

$$(\alpha^1, \alpha^2) \leftrightarrow \begin{pmatrix} (\alpha^1)^3 & 3(\alpha^1)^2 \alpha^2 & 3\alpha^1 (\alpha^2)^2 & (\alpha^2)^3 \\ 0 & (\alpha^1)^2 & 2\alpha^1 \alpha^2 & (\alpha^2)^2 \\ 0 & 0 & \alpha^1 & \alpha^2 \\ 0 & 0 & 0 & 1 \end{pmatrix} \quad (35)$$

DEFINITION. A space is said to be connected if any two points in the space can be joined by a line and all the points of the line lie in the space.

THEOREM. *The component of a continuous group that is connected (called a sheet) with the identity is a group.*

A connected space is simply connected if a curve connecting any two points in the space can be continuously deformed into every other curve connecting the same two points.

DEFINITION. *A Lie group is a group \mathcal{T} which is also an analytic manifold such that the composition function $\phi: \mathcal{T} \times \mathcal{T} \rightarrow \mathcal{T}$ is analytic on its domain of definition [4].*

THEOREM. *The connected component G_0 of a continuous group G is an invariant subgroup of G .*

Example. Transformations on d -dimensional spaces have the general form

$$x'_i = f_i(x_1, x_2, \dots, x_d; a_1, a_2, \dots, a_r), \quad i = 1, 2, \dots, d \quad (36)$$

If the f_i are analytic, then this defines an r -parameter Lie group of transformations.

4.2. Matrix Lie Groups

DEFINITION. *Supposing G is a closed subgroup of $GL(n; \mathbb{C})$, the Lie algebra of G is*

$$\text{Lie}(G) = \{x \in M(n; \mathbb{C}) | e^{tX} \in G \quad \forall t \in \mathbb{R}\} \quad (37)$$

Example. Define the mapping that preserves the operations of the field of complex numbers by

$$\Gamma(1) = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad \Gamma(i) = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \quad (38)$$

then the multiplication structure for the complex numbers

$$\begin{array}{ccc} (a+bi)(c+di) & = & (ac-bd) + (ad+bc)i \\ \downarrow \Gamma & & \uparrow \Gamma^{-1} \\ \begin{pmatrix} a & b \\ -b & a \end{pmatrix} \begin{pmatrix} c & d \\ -d & c \end{pmatrix} & = & \begin{pmatrix} ac-bd & ad+bc \\ -ad-bc & ac-bd \end{pmatrix} \end{array} \quad (39)$$

4.3. Linear Algebras

DEFINITION. *A linear algebra is a vector space V with additional vector multiplication operation:*

- *Closure (Totality):* $\forall a, b \in V, ab \in V$
- *Bilinearity:* $\forall a, b, c \in V, (a+b)c = ac + bc; \quad a(b+c) = ab + ac$

Different kinds of algebras may be obtained, depending on which additional postulates satisfied:

- *Associativity:* $\forall a, b, c \in V, (ab)c = a(bc)$
- *Identity:* $1 \in V$
- *Symmetric/antisymmetric:* $\forall a, b \in V, ab = \pm ba$
- *Derivative property:* $\forall a, b, c \in V, a(bc) = (ab)c + b(ac)$

4.4. Lie Algebras

DEFINITION. *The space of vectors tangent to G at the identity matrix I is called the Lie algebra of the group.*

Given $S \subset \mathbb{R}^k$, a vector v is said to be tangent to S at a point x if there is a differentiable path $\varphi(0) = x$ and $\varphi'(0) = v$. If S is the locus of zeros of one or more polynomial functions $f(x_1, x_2, \dots, x_k)$, it's called a real algebraic set: $S = \{x | f(x) = 0\}$.

Example. the unit circle in \mathbb{R}^2 is a real algebraic set because it's the locus of zeros of the polynomial $f(x_1, x_2) = x_1^2 + x_2^2 - 1 = 0$.

Let $\varphi(t)$ be a path in a real algebraic set S , and let $x = \varphi(t)$, and $v = \varphi'(t)$, thus

$$f(x) = f(\varphi(t)) = 0 \quad (40)$$

Take derivative on both sides

$$0 = \frac{d}{dt} f(\varphi(t)) = \frac{\partial f}{\partial x_1} v_1 + \dots + \frac{\partial f}{\partial x_k} v_k = \nabla f(x) \cdot v \quad (41)$$

where $\nabla f(x) = \left(\frac{\partial f}{\partial x_1}, \dots, \frac{\partial f}{\partial x_k} \right)$ is the gradient vector.

COROLLARY. Let S be a real algebraic set in \mathbb{R}^k , which is the locus of zeros of $f(x)$. The tangent vectors to S at x are orthogonal to the gradients $\nabla f(x)$.

Introduce a formal infinitesimal element ϵ such that $\epsilon^2 = 0$. Define a multiplication on the vector space $E = \{a + b\epsilon \mid a, b \in \mathbb{R}\}$ as the rule

$$\begin{array}{ccc} (a + b\epsilon)(c + d\epsilon) & = & ac + (ad + bc)\epsilon \\ \downarrow \Gamma & & \uparrow \Gamma^{-1} \\ \begin{pmatrix} a & b \\ 0 & a \end{pmatrix} \begin{pmatrix} c & d \\ 0 & c \end{pmatrix} & = & \begin{pmatrix} ac & ad + bc \\ 0 & ac \end{pmatrix} \end{array} \quad (42)$$

where Γ denotes the mapping $\Gamma(1) = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$, $\Gamma(\epsilon) = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$. The addition is vector addition. The main difference between \mathbb{C} and E is that E is not a field, because ϵ has no multiplicative inverse. Instead E is a ring. Since $\epsilon^2 = 0$, the terms of degree ≥ 2 in ϵ drop out. The Taylor's expansion is

$$f(x + v\epsilon) = f(x) + (\nabla f(x) \cdot v)\epsilon \quad (43)$$

DEFINITION. A vector v is called an infinitesimal tangent to a real algebraic set S at x if

$$f(x + v\epsilon) = 0 \quad (44)$$

COROLLARY. Let x be a point of a real algebraic set of S . Every tangent to S at x is an infinitesimal tangent. The converse is also true iff sets S are sufficiently smooth.

Case 1. For $SL_n(\mathbb{R})$, A is an infinitesimal tangent vector if $\det(I + A\epsilon) = 1$.

$$\det(I + A\epsilon) = 1 + (\text{trace}(A))\epsilon = 1 \quad \Rightarrow \quad \text{trace}(A) = 0 \quad (45)$$

PROPOSITION. The following conditions on a real $n \times n$ matrix A are equivalent:

- i. $\text{trace}(A) = 0$;
- ii. e^{tA} is a one-parameter subgroup of $SL_n(\mathbb{R})$;
- iii. A is in the Lie algebra of $SL_n(\mathbb{R})$;
- iv. A is an infinitesimal tangent to $SL_n(\mathbb{R})$ at I .

Case 2. For $O_n(\mathbb{R})$, A is an infinitesimal tangent vector if $(I + A\epsilon)^T(I + A\epsilon) = I$.

$$(I + A\epsilon)^T(I + A\epsilon) = I + (A^T + A)\epsilon = I \quad \Rightarrow \quad A^T + A = 0 \quad (46)$$

PROPOSITION. The following conditions on a real $n \times n$ matrix A are equivalent:

- i. A is skew-symmetric;
- ii. e^{tA} is a one-parameter subgroup of $O_n(\mathbb{R})$;
- iii. A is in the Lie algebra of $O_n(\mathbb{R})$;

iv. A is an infinitesimal tangent to $O_n(\mathbb{R})$ at I .

DEFINITION. The Lie bracket is the law of composition defined by the rule

$$[A, B] = AB - BA \quad (47)$$

The bracket operation is the infinitesimal version of the commutator $PQP^{-1}Q^{-1}$. Using two infinitesimals ϵ, δ and the rules $\epsilon^2 = \delta^2 = 0$ and $\epsilon\delta = \delta\epsilon$. Note $(I + A\epsilon)^{-1} = I - A\epsilon$. If $P = I + A\epsilon$ and $Q = I + B\delta$, the commutator expands to

$$(I + A\epsilon)(I + B\delta)((I + B\delta)(I + A\epsilon))^{-1} = I + (AB - BA)\epsilon\delta \quad (48)$$

Proof. To show the Lie bracket is a law of composition on the Lie algebra, we must check $\forall A, B \in \text{Lie}(G), [A, B] \in \text{Lie}(G)$.

- i. For $G = SL_n(\mathbb{R})$, if $\text{trace}(A) = \text{trace}(B) = 0$, then $\text{trace}([A, B]) = \text{trace}(AB - BA) = \text{trace}(AB) - \text{trace}(BA) = 0$;
- ii. For $G = O_n(\mathbb{R})$, if $A^T = -A, B^T = -B$, then $[A, B]^T = (AB - BA)^T = B^T A^T - A^T B^T = BA - AB = -[A, B]$; \square

Example. If α and β are elements in an abelian group, then $\alpha\beta\alpha^{-1} = \beta$. However, if the group is not commutative, define γ to measure the difference $\alpha\beta\alpha^{-1} = \gamma\beta$. Then γ is a group element because $\alpha\beta(\beta\alpha)^{-1} = \gamma$, which is called the commutator of elements α, β in a group.

Now assuming α, β are close to the identity, expand them in terms of infinitesimal generators:

$$\begin{aligned} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} &= I + \begin{pmatrix} X_\mu \epsilon \\ X_\nu \delta \end{pmatrix} \\ \alpha\beta(\beta\alpha)^{-1} &= I + (X_\mu X_\nu - X_\nu X_\mu)\epsilon\delta \\ &= I + [X_\mu, X_\nu]\epsilon\delta \end{aligned}$$

Since $\alpha\beta(\beta\alpha)^{-1}$ is a group element, the commutator can be expanded in terms of another bases X_λ :

$$[X_\mu, X_\nu] = C_{\mu\nu}^\lambda X_\lambda \quad (49)$$

where $C_{\mu\nu}^\lambda$ are called structure constants.

DEFINITION. A Lie algebra V over a field F is a vector space together with a law of compositions

$$\begin{aligned} V \times V &\rightarrow V \\ v, w &\rightsquigarrow [v, w] \end{aligned} \quad (50)$$

called the Lie bracket, having these properties:

- i. *bilinearity*: $[v_1 + v_2, w] = [v_1, w] + [v_2, w], [v, w_1 + w_2] = [v, w_1] + [v, w_2], [cv, w] = c[v, w] = [v, cw]$;
- ii. *skew symmetry*: $[v, w] = -[w, v]$, or $[v, v] = 0$;
- iii. *Jacobi identity*: $[u, [v, w]] + [v, [w, u]] + [w, [u, v]] = 0$;

for all $u, v, w \in V$ and all $c \in F$.

4.5. Infinitesimal Generators

Let (\mathcal{T}, ϕ) be a Lie group that acts on the geometric space G_N by means of a transformation of coordinates $f(\alpha; x)$, which is a Lie group of transformation. If $F(p)$ is any function defined on all points $p \in G$, then we have in different coordinate system X and Y ,

$$\begin{aligned} F(p) &= F^Y(y^1(p), y^2(p), \dots, y^j(p), \dots, y^N(p)) \\ &= F^X(x^1(p), x^2(p), \dots, x^j(p), \dots, x^N(p)) \end{aligned}$$

The coordinate systems X and Y in G_N are related by a given Lie group of transformation

$$y^j(p) = f^j(\alpha; x(p)) \quad (51)$$

Subsequently we have

$$\begin{aligned} x^j(p) &= f^j(\alpha^{-1}; f(\alpha; x(p))) \\ &= f^j(\alpha^{-1}; y(p)) \\ F^Y(y^1(p), y^2(p), \dots, y^N(p)) &= F^X(f^1(\alpha^{-1}; y(p)), f^2(\alpha^{-1}; y(p)), \dots, f^N(\alpha^{-1}; y(p))) \end{aligned} \quad (52)$$

Concentrate on transformations close to the identity 0, which adds disturbance $\delta\alpha^\mu = \epsilon \rightarrow 0$, $(\delta\alpha^{-1})^\mu = -\epsilon$,

$$\begin{aligned} x^j(p) &= f^j((0 + \delta\alpha^\mu)^{-1}; y(p)) \\ &= f^j(-\epsilon; y(p)) \\ &= f^j(0; y(p)) + \left. \frac{\partial f^j(\alpha; y(p))}{\partial \alpha^\mu} \right|_{\alpha=0} (-\epsilon) + \dots \\ &\cong y^j(p) - \left. \frac{\partial f^j(\alpha; y(p))}{\partial \alpha^\mu} \right|_{\alpha=0} \epsilon \end{aligned} \quad (53)$$

Substitution into (52) generates

$$\begin{aligned} F^Y(y(p)) &= F^X\left(y^j(p) - \left. \frac{\partial f^j(\alpha; y(p))}{\partial \alpha^\mu} \right|_{\alpha=0} \epsilon\right) \\ &\cong F^X(y(p)) - \epsilon \left. \frac{\partial f^j(\alpha; y(p))}{\partial \alpha^\mu} \right|_{\alpha=0} \frac{\partial F^X(y(p))}{\partial y^j} \end{aligned} \quad (54)$$

$$\begin{aligned} F^Y(y) - F^X(y) &= \epsilon \left(- \left. \frac{\partial f^j(\alpha; y)}{\partial \alpha^\mu} \right|_{\alpha=0} \frac{\partial}{\partial y^j} \right) F^X(y) \\ &\equiv \epsilon X_\mu(y) F^X(y) \end{aligned} \quad (55)$$

By changing variables $x = y$,

$$X_\mu(x) = - \left. \frac{\partial f^j(\alpha; x)}{\partial \alpha^\mu} \right|_{\alpha=0} \frac{\partial}{\partial x^j} \quad (56)$$

are defined as the generators of infinitesimal displacements of coordinate systems by ϵ , or simply generators.

Example. For two-parameter group $f^1(\alpha^1, \alpha^2; x) = e^{\alpha^1}x + \alpha^2$, the generators are

$$\begin{aligned} X_1(x) &= - \left. \frac{\partial f(\alpha; x)}{\partial \alpha^1} \right|_{\alpha=0} \frac{\partial}{\partial x} = -x \frac{\partial}{\partial x} \\ X_2(x) &= - \left. \frac{\partial f(\alpha; x)}{\partial \alpha^2} \right|_{\alpha=0} \frac{\partial}{\partial x} = - \frac{\partial}{\partial x} \end{aligned} \quad (57)$$

Suppose we have a function $F^X(x) = (x - c)^2$,

Case 1. If there exists a relation $y = x + \epsilon$, where $\epsilon = \delta\alpha^\mu$, $\mu = 2$, then

$$\begin{aligned} F^Y(y) &= (y - \epsilon - c)^2 \\ &= (y - c)^2 + \epsilon X_\mu(y) F^X(y) \\ &= (I + \epsilon X_\mu(y))(y - c)^2 \end{aligned}$$

for finite displacement $y = x + \alpha^\mu$, $\mu = 2$,

$$\begin{aligned} F^Y(y) &= \lim_{N \rightarrow \infty} \left(I + \frac{\alpha^2}{N} X_2(y) \right)^N (y - c)^2 \\ &= e^{\alpha^2 X_2(y)} (y - c)^2 \\ &= e^{-\alpha^2 \frac{\partial}{\partial y}} (y - c)^2 \\ &= (y - (c + \alpha^2))^2 \end{aligned}$$

Case 2. If there exists a relation $y = e^{\delta\alpha^1}x = (1 + \delta\alpha^1)x$, then

$$F^Y(y) = (I - \delta\alpha^1 X_1(y))(y - c)^2$$

for finite displacement $y = x + \alpha^\mu$, $\mu = 1$,

$$\begin{aligned} F^Y(y) &= \lim_{N \rightarrow \infty} \left(I - \frac{\alpha^1}{N} X_1(y) \right)^N (y - c)^2 \\ &= e^{-\alpha^1 X_2(y)} (y - c)^2 \\ &= e^{-\alpha^1 y \frac{\partial}{\partial y}} (y - c)^2 \\ &= (e^{-\alpha^1} y - c)^2 \end{aligned}$$

The idea [5] behind infinitesimal generator is to consider an infinitesimal transformation around the identity. Any finite transformation can be then constructed by the integration of this infinitesimal transformation.

Example. For SO_2 , expand $R(\theta) = \begin{pmatrix} \cos(\theta) & -\sin(\theta) \\ \sin(\theta) & \cos(\theta) \end{pmatrix}$ into a Taylor series around the identity ($\theta = 0$):

$$R(\theta) = R(0) + \frac{dR}{d\theta} \Big|_{\theta=0} \theta + \frac{1}{2} \frac{d^2 R}{d\theta^2} \Big|_{\theta=0} \theta^2 + \dots \quad (58)$$

From $R(\theta_1 + \theta_2) = R(\theta_1)R(\theta_2)$, we have

$$\frac{dR(\theta_1 + \theta_2)}{d\theta_1} = \frac{dR(\theta_1)}{d\theta_1} R(\theta_2) \quad (59)$$

With the setting $\theta_1 = 0$, the left-hand side of (59) is

$$\frac{dR(\theta_1 + \theta_2)}{d(\theta_1 + \theta_2)} \frac{d(\theta_1 + \theta_2)}{d\theta_1} \Big|_{\theta_1=0} = \frac{dR(\theta_1 + \theta_2)}{d(\theta_1 + \theta_2)} \Big|_{\theta_1=0} = \frac{dR(\theta_2)}{d\theta_2} \quad (60)$$

the right-hand side of (59) is

$$\frac{dR(\theta_1)}{d\theta_1} R(\theta_2) \Big|_{\theta_1=0} = \begin{pmatrix} -\sin(\theta_1) & -\cos(\theta_1) \\ \cos(\theta_1) & -\sin(\theta_1) \end{pmatrix} \Big|_{\theta_1=0} R(\theta_2) = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} R(\theta_2) \quad (61)$$

then after changing variable (59) becomes

$$\frac{dR(\theta)}{d\theta} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} R(\theta) \equiv X R(\theta) \quad (62)$$

By observing that $R(0) = I$, we obtain

$$\frac{dR(\theta)}{d\theta} \Big|_{\theta=0} = X R(0) = X \quad (63)$$

The higher-order derivatives of R is then

$$\frac{d^n R(\theta)}{d\theta^n} \Big|_{\theta=0} = X \frac{d^{n-1} R(\theta)}{d\theta^{n-1}} \Big|_{\theta=0} = X^2 \frac{d^{n-2} R(\theta)}{d\theta^{n-2}} \Big|_{\theta=0} = \dots = X^n \quad (64)$$

The Taylor series becomes

$$\begin{aligned} R(\theta) &= I + X\theta + \frac{1}{2} X^2 \theta^2 + \dots \\ &= \sum_{n=0}^{\infty} \frac{1}{n!} (X\theta)^n \\ &= e^{\theta X} \end{aligned} \quad (65)$$

$$e^{\theta X} = I \cos(\theta) + X \sin(\theta) \quad (66)$$

Thus every rotation by a finite angle can be obtained from the exponentiation of the matrix X , which is called the infinitesimal generator of rotations.

Example. $GL_3(\mathbb{R})$ has 9 parameters, but the invariance of the length produces 6 independent conditions, leaving 3 free parameters, so $O_3(\mathbb{R})$ forms a 3-parameter Lie group.

1. First rotation is about the z -axis by angle θ_3 :

$$R_3(\theta_3) = \begin{pmatrix} \cos(\theta_3) & -\sin(\theta_3) & 0 \\ \sin(\theta_3) & \cos(\theta_3) & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad X_3 = \left. \frac{dR_3}{d\theta_3} \right|_{\theta_3=0} = \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad (67)$$

2. Second rotation is about the x -axis by angle θ_1 :

$$R_1(\theta_1) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos(\theta_1) & -\sin(\theta_1) \\ 0 & \sin(\theta_1) & \cos(\theta_1) \end{pmatrix} \quad X_1 = \left. \frac{dR_1}{d\theta_1} \right|_{\theta_1=0} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} \quad (68)$$

3. Final rotation is about the y -axis by angle θ_2 :

$$R_2(\theta_2) = \begin{pmatrix} \cos(\theta_2) & 0 & \sin(\theta_2) \\ 0 & 1 & 0 \\ -\sin(\theta_2) & 0 & \cos(\theta_2) \end{pmatrix} \quad X_2 = \left. \frac{dR_2}{d\theta_2} \right|_{\theta_2=0} = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{pmatrix} \quad (69)$$

The rotations do not commute. Define the commutator of X_1 and X_2 as

$$[X_1, X_2] \equiv X_1 X_2 - X_2 X_1 = X_3 \quad [X_2, X_3] = X_1 \quad [X_3, X_1] = X_2 \quad (70)$$

Another approach is to make small disturbance

$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} 1 & -\theta_3 & \theta_2 \\ \theta_3 & 1 & -\theta_1 \\ -\theta_2 & \theta_1 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

Substituting this coordinate transformation into a differentiable function $F(x, y, z)$,

$$\begin{aligned} F(x', y', z') &= F(x - \theta_3 y + \theta_2 z, \theta_3 x + y - \theta_1 z, -\theta_2 x + \theta_1 y + z) \\ &= F(x, y, z) + \left(\frac{\partial F}{\partial z} y - \frac{\partial F}{\partial y} z \right) \theta_1 + \left(\frac{\partial F}{\partial x} z - \frac{\partial F}{\partial z} x \right) \theta_2 + \left(\frac{\partial F}{\partial y} x - \frac{\partial F}{\partial x} y \right) \theta_3 \end{aligned} \quad (71)$$

The generators X_i are

$$\begin{aligned} X_1 &= y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \\ X_2 &= z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \\ X_3 &= x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \end{aligned} \quad (72)$$

4.5.1. Infinitesimal generators for Lie groups

Instead of geometric space G , consider in topological space \mathcal{T} , in which $\gamma = \phi(\beta, \alpha)$. The infinitesimal generators are

$$\begin{aligned} X_\mu(x) &= - \left. \frac{\partial f^i(\beta; x)}{\partial \beta^\mu} \right|_{\beta=0} \frac{\partial}{\partial x^i} \\ X_\mu(\alpha) &= - \left. \frac{\partial \phi^\lambda(\beta; \alpha)}{\partial \beta^\mu} \right|_{\beta=0} \frac{\partial}{\partial \alpha^\lambda} \end{aligned} \quad (73)$$

A Lie group acting on itself is a nonsingular change of basis,

$$\det \left| \left. \frac{\partial \phi^\lambda(\beta; \alpha)}{\partial \beta^\mu} \right|_{\beta=0} \right| \neq 0$$

Define

$$\mathbf{u}(x) = \left. \frac{\partial f^i(\beta; x)}{\partial \beta^\mu} \right|_{\beta=0} \quad \Psi(\alpha) = \left. \frac{\partial \phi^\lambda(\beta; \alpha)}{\partial \beta^\mu} \right|_{\beta=0} \quad (74)$$

then we have

$$\begin{aligned} X_\mu(x) &= -\mathbf{u}(x) \frac{\partial}{\partial x^i} \\ X_\mu(\alpha) &= -\Psi(\alpha) \frac{\partial}{\partial \alpha^\lambda} \end{aligned} \quad (75)$$

Example. If we define a Lie group and its multiplication law as

$$\begin{aligned} \phi^1(\beta^1 \beta^2; \chi^1 \chi^2) &= \beta^1 + \chi^1 \\ \phi^2(\beta^1 \beta^2; \chi^1 \chi^2) &= e^{\beta^1} \chi^2 + \beta^2 \end{aligned}$$

The generators are

$$\begin{pmatrix} X_1(\chi) \\ X_2(\chi) \end{pmatrix} = - \left(\begin{array}{cc} \frac{\partial \phi^1}{\partial \beta^1} = 1 & \frac{\partial \phi^2}{\partial \beta^1} = \chi^2 \\ \frac{\partial \phi^1}{\partial \beta^2} = 0 & \frac{\partial \phi^2}{\partial \beta^2} = 1 \end{array} \right) \bigg|_{\beta=0} \left(\begin{array}{c} \frac{\partial}{\partial \chi^1} \\ \frac{\partial}{\partial \chi^2} \end{array} \right) = \begin{pmatrix} -\frac{\partial}{\partial \chi^1} - \chi^2 \frac{\partial}{\partial \chi^2} \\ -\frac{\partial}{\partial \chi^2} \end{pmatrix} \quad (76)$$

The group operation can be interpreted in two ways, because the dimensions are equal for two parameters:

1. Left translation by α : $\chi'(p) = \phi(\beta; \chi(p))$; the generators are as in (76).
2. Right translation by χ : $\beta'(p) = \phi(\beta(p); \chi)$; the generators are

$$\begin{pmatrix} X_1(\beta) \\ X_2(\beta) \end{pmatrix} = - \left(\begin{array}{cc} \frac{\partial \phi^1}{\partial \chi^1} = 1 & \frac{\partial \phi^2}{\partial \chi^1} = 0 \\ \frac{\partial \phi^1}{\partial \chi^2} = 0 & \frac{\partial \phi^2}{\partial \chi^2} = e^{\beta^1} \end{array} \right) \bigg|_{\chi=0} \left(\begin{array}{c} \frac{\partial}{\partial \beta^1} \\ \frac{\partial}{\partial \beta^2} \end{array} \right) = \begin{pmatrix} -\frac{\partial}{\partial \beta^1} \\ -e^{\beta^1} \frac{\partial}{\partial \beta^2} \end{pmatrix} \quad (77)$$

4.5.2. Infinitesimal generators for matrix groups

If $M(\alpha^1, \alpha^2, \dots, \alpha^n)$ is an element of a group of $r \times r$ matrices, the infinitesimal generators are

$$X_\mu(r \times r) = \lim_{\alpha^\mu \rightarrow 0} \frac{M(0, 0, \dots, \alpha^\mu, \dots, 0) - M(0, 0, \dots, 0, \dots, 0)}{\alpha^\mu}$$

4.6. Lie's Three Theorems

4.6.1. Lie's First Theorem

THEOREM. If $y^j(p) = f^j(\alpha; x(p))$ is analytic, then

$$\frac{\partial y^j}{\partial \alpha^\mu} = \frac{\partial f^j(\alpha; x)}{\partial \alpha^\mu} = \sum_{k=0}^n \Psi_{\mu k}(\alpha) \mathbf{u}_{kj}(y) \quad (78)$$

Proof. If we transform from coordinate system S to S' , and then to S'' ,

$$S(x) \xrightarrow{\alpha} S'(y) \xrightarrow{\beta = \delta \alpha^\mu} S''(z)$$

where β is close to the identity. Then the difference

$$dy^j = z^j - y^j$$

can be computed as

$$\begin{aligned} y^j &= f^j(\alpha; x) \\ z^j &= (y + \delta y)^j \\ &= f^j(\delta \alpha^\mu; y) \\ dy^j &= \left. \frac{\partial f^j(\alpha; y)}{\partial \alpha^\mu} \right|_{\alpha=0} \delta \alpha^\mu \\ &\equiv \delta \alpha^\mu \mathbf{u}(y) \end{aligned} \quad (79)$$

where $\mathbf{u}(y)$ is a $\eta \times N$ nonsingular matrix. Also we have

$$\begin{aligned} (\alpha + d\alpha)^\lambda &= \phi^\lambda(\delta\alpha; \alpha) \\ d\alpha^\lambda &= \left. \frac{\partial \phi^\lambda(\beta; \alpha)}{\partial \beta^\mu} \right|_{\beta=0} \delta\alpha^\mu \\ &\equiv \delta\alpha^\mu \Theta(\alpha) \end{aligned}$$

Since $\Theta(\alpha)$ is an $\eta \times \eta$ nonsingular matrix, it has an inverse $\Psi(\alpha)$,

$$\begin{aligned} \Psi(\alpha)\Theta(\alpha) &= \Theta(\alpha)\Psi(\alpha) = I \\ \delta\alpha^\mu &= d\alpha^\lambda \Psi(\alpha) \end{aligned}$$

Substitution into (79) generates

$$\begin{aligned} dy^j &= \delta\alpha^\mu \mathbf{u}(y) \\ &= d\alpha^\lambda \Psi(\alpha) \mathbf{u}(y) \\ \frac{\partial y^j}{\partial \alpha^\mu} &= \sum_{k=0}^{\eta} \Psi_{\mu k}(\alpha) \mathbf{u}_{kj}(y) \end{aligned}$$

In this theorem Lie decoupled PDE into the product of two matrices. □

Example. If $\alpha = (\alpha^1, \alpha^2)$ and $\delta\alpha = (\delta\alpha^1, \delta\alpha^2)$ is close to the identity, plus

$$\begin{aligned} y &= e^{\alpha^1} x + \alpha^2 \\ y + dy &= e^{\delta\alpha^1} y + \delta\alpha^2 \\ dy &= \delta\alpha^1 y + \delta\alpha^2 \end{aligned} \tag{80}$$

Since $\delta\alpha$ is close to the identity, we can write

$$(\alpha + d\alpha)^\mu = \phi^\mu(\delta\alpha^1, \delta\alpha^2; \alpha^1, \alpha^2) \tag{81}$$

Because

$$\begin{aligned} f(\beta, f(\alpha; x)) &= f(\beta, e^{\alpha^1} x + \alpha^2) \\ &= e^{\beta^1} (e^{\alpha^1} x + \alpha^2) + \beta^2 \\ &= e^{\beta^1 + \alpha^1} x + (\alpha^2 e^{\beta^1} + \beta^2) \\ &= f(\beta^1 + \alpha^1, \alpha^2 e^{\beta^1} + \beta^2; x) \\ \phi(\beta; \alpha) &= (\beta^1 + \alpha^1, \alpha^2 e^{\beta^1} + \beta^2) \end{aligned}$$

then with $\beta = \delta\alpha$ (81) becomes

$$\begin{aligned} \begin{pmatrix} \alpha^1 + d\alpha^1 \\ \alpha^2 + d\alpha^2 \end{pmatrix} &= \begin{pmatrix} \delta\alpha^1 + \alpha^1 \\ \alpha^2 e^{\delta\alpha^1} + \delta\alpha^2 \end{pmatrix} \\ &= \begin{pmatrix} \delta\alpha^1 + \alpha^1 \\ \alpha^2(1 + \delta\alpha^1) + \delta\alpha^2 \end{pmatrix} \\ \begin{pmatrix} d\alpha^1 \\ d\alpha^2 \end{pmatrix} &= \begin{pmatrix} \delta\alpha^1 \\ \alpha^2 \delta\alpha^1 + \delta\alpha^2 \end{pmatrix} \\ &= \begin{pmatrix} 1 & 0 \\ \alpha^2 & 1 \end{pmatrix} \begin{pmatrix} \delta\alpha^1 \\ \delta\alpha^2 \end{pmatrix} \end{aligned}$$

Take the inverse for both sides

$$\begin{aligned} \begin{pmatrix} \delta\alpha^1 \\ \delta\alpha^2 \end{pmatrix} &= \begin{pmatrix} 1 & 0 \\ \alpha^2 & 1 \end{pmatrix}^{-1} \begin{pmatrix} d\alpha^1 \\ d\alpha^2 \end{pmatrix} \\ &= \begin{pmatrix} 1 & 0 \\ -\alpha^2 & 1 \end{pmatrix} \begin{pmatrix} d\alpha^1 \\ d\alpha^2 \end{pmatrix} \end{aligned}$$

Substitution into (80) produces

$$\begin{aligned} dy &= \delta\alpha^1 y + \delta\alpha^2 \\ &= d\alpha^1 y + (-\alpha^2 d\alpha^1 + d\alpha^2) \\ &= (y - \alpha^2)d\alpha^1 + d\alpha^2 \end{aligned}$$

which means

$$\begin{aligned} \frac{\partial y}{\partial \alpha^1} &= y - \alpha^2 \\ \frac{\partial y}{\partial \alpha^2} &= 1 \end{aligned}$$

The second method to determine these partial derivatives is

$$\begin{aligned} \Theta(\alpha) &= \left. \frac{\partial \phi^\lambda(\beta; \alpha)}{\partial \beta^\mu} \right|_{\beta=0} = \left(\begin{array}{cc} \frac{\partial \phi^1}{\partial \beta^1} & \frac{\partial \phi^2}{\partial \beta^1} \\ \frac{\partial \phi^1}{\partial \beta^2} & \frac{\partial \phi^2}{\partial \beta^2} \end{array} \right)_{\beta=0} = \left(\begin{array}{cc} 1 & \alpha^2 \\ 0 & 1 \end{array} \right) \\ \mathbf{u}(y) &= \left. \frac{\partial f^j(\alpha; y)}{\partial \alpha^\mu} \right|_{\alpha=0} = \left(\begin{array}{c} \frac{\partial f}{\partial \alpha^1} \\ \frac{\partial f}{\partial \alpha^2} \end{array} \right)_{\alpha=0} = \left(\begin{array}{c} y \\ 1 \end{array} \right) \end{aligned}$$

Then we have

$$\frac{\partial y^j}{\partial \alpha^\mu} = \sum_{k=0}^{\eta} \Psi_{\mu k}(\alpha) \mathbf{u}_{kj}(y) = \left(\begin{array}{cc} 1 & \alpha^2 \\ 0 & 1 \end{array} \right)^{-1} \left(\begin{array}{c} y \\ 1 \end{array} \right) = \left(\begin{array}{c} y - \alpha^2 \\ 1 \end{array} \right)$$

4.6.2. Lie's Second Theorem

THEOREM. If X_μ are the generators of a Lie group, then the coefficients $C_{\mu\nu}^\lambda$ given by

$$[X_\mu, X_\nu] = C_{\mu\nu}^\lambda X_\lambda \quad (82)$$

are constants.

DEFINITION. The necessary and sufficient conditions for the existence of a unique solution with initial conditions

$$y^i = f^i(\alpha; x)|_{\alpha=0} = x^i \quad (83)$$

is that all mixed derivatives are equal:

$$\frac{\partial^2 y^i}{\partial \alpha^\mu \partial \alpha^\nu} = \frac{\partial^2 y^i}{\partial \alpha^\nu \partial \alpha^\mu} \quad (84)$$

which is called integrability conditions.

Proof. Apply (84) in Lie's first theorem:

$$\begin{aligned} \frac{\partial}{\partial \alpha^\mu} (\Psi_{\nu k}(\alpha) \mathbf{u}_{kj}(x)) &= \frac{\partial}{\partial \alpha^\nu} (\Psi_{\mu \lambda}(\alpha) \mathbf{u}_{\lambda j}(x)) \\ \Psi_{\nu k}(\alpha) \frac{\partial \mathbf{u}_{kj}(x)}{\partial \alpha^\mu} - \Psi_{\mu \lambda}(\alpha) \frac{\partial \mathbf{u}_{\lambda j}(x)}{\partial \alpha^\nu} &= \frac{\partial \Psi_{\mu \lambda}(\alpha)}{\partial \alpha^\nu} \mathbf{u}_{\lambda j}(x) - \frac{\partial \Psi_{\nu k}(\alpha)}{\partial \alpha^\mu} \mathbf{u}_{kj}(x) \end{aligned} \quad (85)$$

Replace the terms $\frac{\partial \mathbf{u}}{\partial \alpha}$ on the left by

$$\begin{aligned} \frac{\partial \mathbf{u}_{kj}(x)}{\partial \alpha^\mu} &= \frac{\partial x^i}{\partial \alpha^\mu} \frac{\partial \mathbf{u}_{kj}(x)}{\partial x^i} = \Psi_{\mu k}(\alpha) \mathbf{u}_{kj}(x) \frac{\partial \mathbf{u}_{kj}(x)}{\partial x^i} \\ \frac{\partial \mathbf{u}_{\lambda j}(x)}{\partial \alpha^\nu} &= \frac{\partial x^i}{\partial \alpha^\nu} \frac{\partial \mathbf{u}_{\lambda j}(x)}{\partial x^i} = \Psi_{\nu \lambda}(\alpha) \mathbf{u}_{\lambda j}(x) \frac{\partial \mathbf{u}_{\lambda j}(x)}{\partial x^i} \end{aligned} \quad (86)$$

Equation (85) then becomes

□

4.6.3. Lie's Third Theorem

THEOREM. *The structure constants obey*

$$C_{\mu\nu}^{\lambda} = -C_{\nu\mu}^{\lambda} \quad (87)$$

$$C_{\alpha\beta}^{\delta}C_{\delta\gamma}^{\rho} + C_{\beta\gamma}^{\delta}C_{\delta\alpha}^{\rho} + C_{\gamma\alpha}^{\delta}C_{\delta\beta}^{\rho} = 0 \quad (88)$$

Proof. We have

$$C_{\mu\nu}^{\lambda} = \left(\frac{\partial^2 \phi^{\lambda}(\beta; \alpha)}{\partial \beta^{\mu} \partial \alpha^{\nu}} - \frac{\partial^2 \phi^{\lambda}(\beta; \alpha)}{\partial \beta^{\nu} \partial \alpha^{\mu}} \right) \Big|_{\beta=\alpha=0} = -C_{\nu\mu}^{\lambda}$$

Equation (88) is a trivial consequence of the Jacobi identity.

$$[[X_{\alpha}, X_{\beta}], X_{\gamma}] + [[X_{\beta}, X_{\gamma}], X_{\alpha}] + [[X_{\gamma}, X_{\alpha}], X_{\beta}] = 0 \quad (89)$$

The Jacobi identity (89) bears a strong resemblance to the equation

$$\frac{d}{dx}(f(x)g(x)) = \frac{df(x)}{dx}g(x) + f(x)\frac{dg(x)}{dx} \quad (90)$$

For this reason, the Lie bracket $[\cdot, \cdot]$ is sometimes called a derivative. □

4.7. Converse of Lie's Three Theorems

4.8. Taylor's Theorem for Lie Groups

THEOREM. *There exists an analytic mapping $\gamma^{\mu} = \phi^{\mu}(\beta, \alpha)$ in which every straight line through the origin is a one-dimensional abelian subgroup. The Lie group operation corresponding to the Lie algebra element $\alpha^{\mu}X_{\mu}$ is*

$$\alpha^{\mu}X_{\mu} \rightarrow e^{-\alpha^{\mu}X_{\mu}}$$

Proof. Since $X_{\mu}(x) = -\mathbf{u}(x)\frac{\partial}{\partial x^i}$, we can write

$$\frac{\partial x^j}{\partial \alpha^{\mu}} = \Psi(\alpha)\mathbf{u}(x) = -\Psi(\alpha)X_{\mu}(x)x^i$$

When we look at the straight line

$$\alpha^{\mu}(\tau) = s^{\mu}\tau$$

through the origin of the Lie algebra, the x^i are functions of the single parameter τ :

$$\frac{dx^i(\tau)}{d\tau} = \frac{\partial x^i}{\partial \alpha^{\lambda}} \frac{d\alpha^{\lambda}}{d\tau} = -\Psi(\alpha = s\tau)X_{\mu}(x)x^i s^{\lambda}$$

□

4.9. Classical Matrix Groups

4.9.1. Quaternion Groups

An arbitrary quaternion in $GL(1, q)$ is written

$$q = q_0\lambda_0 + q_1\lambda_1 + q_2\lambda_2 + q_3\lambda_3 \quad (91)$$

in which λ_i obey the multiplication

	λ_0	λ_1	λ_2	λ_3
λ_0	λ_0	λ_1	λ_2	λ_3
λ_1	λ_1	$-\lambda_0$	λ_3	$-\lambda_2$
λ_2	λ_2	$-\lambda_3$	$-\lambda_0$	λ_1
λ_3	λ_3	λ_2	$-\lambda_1$	$-\lambda_0$

Table 1. Quaternion bases

- i. The identity under multiplication is $1\lambda_0$;
- ii. The identity under addition is 0;
- iii. The quaternions form a group under multiplication excluding 0, which is a field.

4.9.2. Unitary Groups

4.9.3. Orthogonal Groups

4.10. Lie's Theory on PDE

If λ_0 indicates the identity transformation,

$$x = X(x, y, \lambda_0) \quad y = Y(x, y, \lambda_0) \quad (92)$$

expand it in a Taylor's series around $\lambda = \lambda_0$ for sufficiently small $\lambda - \lambda_0$,

$$\begin{aligned} \begin{pmatrix} x' \\ y' \end{pmatrix} &= \begin{pmatrix} x \\ y \end{pmatrix} + \begin{pmatrix} \frac{\partial X}{\partial \lambda} \\ \frac{\partial Y}{\partial \lambda} \end{pmatrix}_{\lambda=\lambda_0} (\lambda - \lambda_0) \\ &\equiv \begin{pmatrix} x \\ y \end{pmatrix} + \begin{pmatrix} \xi(x, y) \\ \eta(x, y) \end{pmatrix} (\lambda - \lambda_0) \end{aligned} \quad (93)$$

The infinitesimal transformation (93) is Euler's finite-difference algorithm for solving the coupled differential equations

$$\frac{dx}{\xi(x, y)} = \frac{dy}{\eta(x, y)} = d\lambda \quad (94)$$

Example. For $\lambda_0 = 0$ and $\xi(x, y) = -y, \eta(x, y) = x$. Thus (94) becomes

$$\begin{aligned} \frac{dx}{-y} &= \frac{dy}{x} = d\lambda \\ x dx &= -y dy \\ \frac{1}{2}x^2 &= -\frac{1}{2}y^2 + \frac{1}{2}a^2 \\ x^2 + y^2 &= a^2 \end{aligned} \quad (95)$$

The second equation of (95) can be written

$$\begin{aligned} \frac{dy}{\sqrt{a^2 - y^2}} &= d\lambda \\ \arcsin\left(\frac{y}{a}\right) &= \lambda + b \\ y &= a \sin(\lambda + b) \\ x &= a \cos(\lambda + b) \end{aligned}$$

5. INTEGRATORS

5.1. Euler Method

- Explicit Euler Method [3]

$$y_{n+1} = y_n + hf(y_n) \quad (96)$$

- Implicit Euler Method

$$y_{n+1} = y_n + hf(y_{n+1}) \quad (97)$$

- Implicit Midpoint Rule

$$y_{n+1} = y_n + hf\left(\frac{y_n + y_{n+1}}{2}\right) \quad (98)$$

- Symplectic Euler Method
For partitioned systems

$$\begin{aligned} \dot{u} &= a(u, v) \\ \dot{v} &= b(u, v) \end{aligned} \quad (99)$$

partitioned Euler methods which treats one variable by the implicit and the other variable by the explicit Euler method

$$\begin{aligned} u_{n+1} &= u_n + ha(u_n, v_{n+1}) \\ v_{n+1} &= v_n + hb(u_n, v_{n+1}) \end{aligned} \quad \text{or} \quad \begin{aligned} u_{n+1} &= u_n + ha(u_{n+1}, v_n) \\ v_{n+1} &= v_n + hb(u_{n+1}, v_n) \end{aligned} \quad (100)$$

is a symplectic Euler method.

5.2. Symplectic Integrators

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